Microstructure by Design: Integrating Grain Growth Experiments, Data Analytics, Simulation and Theory



Katayun Barmak¹, Yekaterina Epshteyn², Chun Liu³, Jeffrey Rickman⁴ ¹Department of Applied Physics and Applied Mathematics, Columbia University, New York, NY 10027; ²Department of Mathematics, University of Utah, Salt Lake City, UT 84112; ³Department of Applied Mathematics, Illinois Institute of Technology, Chicago, IL 60616; ⁴Department of Materials Science and Engineering, Lehigh University, Bethlehem, PA 18015

Project Scope

The principal objective of this project is to quantify statistical measures of grain growth obtained from experimental data garnered from metallic films and to employ these measures to inform mesoscale models of grain growth. This link forged between experiment and theory via analysis and simulations will result in improved models of grain coarsening that properly reflect the underlying physics. Data analytics will be employed to study the experimental evolution computational Of and microstructures, regarded here as a collection of interacting grain triple junctions, to validate and refine the models and guide future experiments

Relevance to MGI

The MGI goal of rapid discovery and deployment of advanced materials requires, in the case of polycrystalline materials, an enhanced understanding of microstructural development to guide subsequent materials design. Thus, this project integrates experiment, theory and simulation to obtain a better, quantitative understanding of grain growth mechanisms, using thin metallic films as the experimental testbed.

Technical Progress

- Achieved automated grain boundary detection in bright-field TEM images using a modified U-Net convolutional neural network.
- Experimentally obtained correlations between grain boundary character distribution (GBCD) and grain boundary energy distribution (GBED) imply that the force balance at the triple junction (Herring condition) does not fully specify triple junction geometry.
- Developed statistical tools to distribution of triple junctions in experimental crystal orientation maps.

References

(1) Y. Epshteyn, C. Liu, M. Mizuno, SIMA, 53(3), 3072-3097, 2021. (2) K. Barmak, A. Dunca, Y. Epshteyn, C. Liu, M. Mizuno, AWM-Springer Volume, In Press, arXiv:2105.07255, 2022. (3) Y. Epshteyn, C. Liu, M. Mizuno, Submitted, arXiv:2106.14249, 2021.

characterize

Established thermodynamically consistent models for deterministic evolution, as well as stochastic dynamics of grain boundaries. The new models incorporate and explore the specific dynamics of misorientations and the triple junctions, in contrast to the conventional approaches with equilibrium Herring conditions. Systematically studied the mathematical structures of the derived systems. Designed corresponding structure-preserving numerical validate the analytical results, as well as to validate against current and future experimental data.



Distribution of Triple Junctions in a Tungsten Film

algorithms to

Different Time Scales of 2D Coarsening Networks: Curvature, Mobility of Triple Junctions and Misorientations Dynamics

Total grain boundary energy: $E(t) = \sum_{r_k} \psi^k (\Delta \alpha^k) d\mathcal{H}^k$

$$\frac{\mathrm{d}E(t)}{\mathrm{d}t} = -\sum_{k} \int_{\Gamma^{k}} \frac{1}{\mu^{k}} \left| v_{n}^{k} \right|^{2} \mathrm{d}\mathscr{H}^{1} - \sum_{m} \frac{1}{\eta^{m}} \left| v^{m} \right|^{2} - \sum_{l} \frac{1}{\gamma^{k}} \left| \frac{\mathrm{d}\alpha^{k}}{\mathrm{d}t} \right|^{2}$$

Grain boundary energy density:

 $\psi(\Delta \alpha) = 1 + 0.25 \sin^2(2\Delta \alpha)$

Motion of interfaces: $v_n = \mu \psi \kappa$ on Γ Dynamics of misorientations: $\frac{d(\Delta \alpha)}{dt} = -\gamma \int_{\Gamma} \frac{\partial \psi}{\partial (\Delta \alpha)} \ d\mathcal{H}^{1}$ Motion of triple junctions: $v^m = -\eta^m \sum T^{m,\ell}$

The associated probability density function $f(\Delta \alpha, a, t)$ obeys the Fokker-Planck

Effect of dynamic misorientations relaxation time scale on Grain Boundary Character Distribution (GBCD) Marginal probability density: $\rho_1(\Delta \alpha, t) = \int_{\Omega_{TI}} f(\Delta \alpha, a, t) da$; the Herring angle condition $(\eta \to \infty)$

Platinum film from instance of an in-situ grain growth experiment

equation $\frac{\partial f}{\partial t} + \nabla_{\Delta \alpha}^{\Omega} \cdot (v_{\Delta \alpha} f) + \nabla_a \cdot (v_a f) = \frac{\beta_{\Delta \alpha}^2}{2} \Delta_{\Delta \alpha}^{\Omega} f + \frac{\beta_a^2}{2} \Delta_a f$ subject to the natural boundary condition. If the relaxation time scales and the fluctuation parameters satisfy $\frac{6\gamma}{R^2} = \frac{2\eta}{R^2}$, the Fokker-Planck equation has the energy law $(D := \frac{\beta_{\Delta \alpha}^2}{6\gamma})$:

 $\frac{\mathrm{d}}{\mathrm{d}t} \iint_{\Omega \times \Omega_{\mathrm{TI}}} (Df \log f + fE) \mathrm{d}\Delta\alpha \mathrm{d}a = -\frac{\beta_{\Delta\alpha}^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_{\Delta\alpha}^{\Omega} (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI}}} f \left| \nabla_a (D \log f + E) \right|^2 \mathrm{d}\Delta\alpha \mathrm{d}a - \frac{\beta_a^2}{2D} \iint_{\Omega \times \Omega_{\mathrm{TI$

principle

 Boltzmann distribu Misorientation

- Established and verified in grain growth simulations sufficient condition (which connects triple junctions geometry with grain boundary energies) to observe

- Consistent with the dissipation-fluctuation

Boltzmann distribution for the grain boundary energy density as a steady-state distribution for the GBCD.

> Instance of the simulation: **No-curvature**